

# Exploring ULXs as Short GRB Precursors

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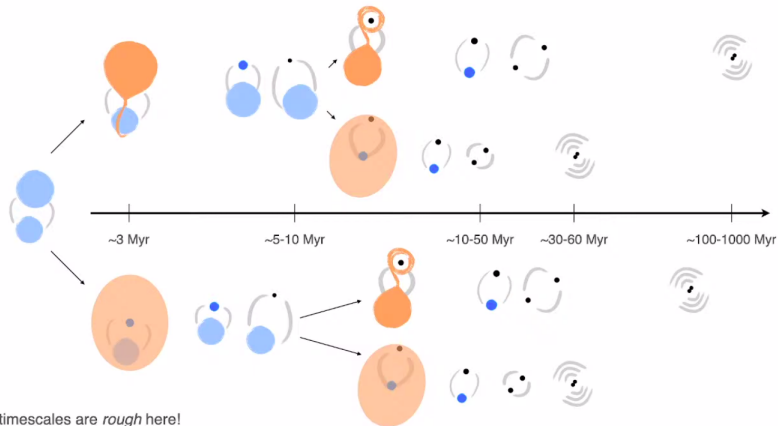


**CAPP**  
Centre for Astro-Particle Physics

- Stellar evolution
- Research Objectives
- Method and ULX model
- Results
- Summary

# Stellar Evolution

## DCO formation channels:



# What are ULXs

Ultra-Luminous X-ray sources (ULXs) are galactic point-like sources, exhibiting high X-ray luminosity ( $L_x > 10^{39}$  erg/s) exceeding the Eddington limit of neutron stars or stellar-mass black holes ([Fabbiano 1989](#)).

- Early surveys identified 16 ULXs with  $L_x > 10^{39}$  erg/s ([Fabbiano 1989](#)).
- Over 1800 ULXs identified so far ([Walton et al. 2022](#)).
- ULXs mainly in extragalactic regions, not in nuclei, are distinct from accreting supermassive black holes.

# Research Objectives

- Simulate binary systems ( using a population synthetic code)
- Predict DCO merger rates from ULX phase evolution.
- Predict SGRB rates from DNS that evolved through the ULX phase.

- **COSMIC (Compact Object Synthesis and Monte Carlo Investigation Code) developed by Breivik et al. (2020):**
  - Population synthesis code used for simulating the evolution of binary systems.
  - Combines theoretical models of stellar evolution, binary interactions, and compact object mergers.
- **Initial Set-up:**
  - Primary star mass ( $M_1$ ) drawn from a Kroupa initial mass function (IMF) (Kroupa et al. 1993).
  - Secondary star mass ( $M_2$ ) drawn from a flat distribution of the binary mass ratio ( $q = M_2/M_1$ ).
- **Simulation Details:**
  - Simulated  $2 \times 10^6$  binary systems at different metallicities (0.5% to 150%  $Z_\odot$ ).
  - Evolved binaries from ZAMS to X-ray binary phase and then to the formation of DCO.
  - Allowed super-Eddington accretion up to 1000 times Eddington rate.

# Disc Model

We adopt the SCAD model ( [Vinokurov et al. 2013](#)).

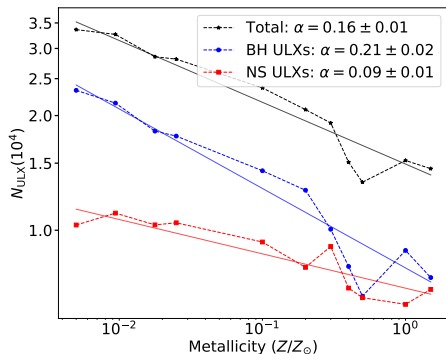
$$T(r) = T_{\text{in}} \begin{cases} (r \sin \theta_f)^{-1/2}, & 1 \leq r \leq r_{\text{sp}} \\ \left[ \frac{f_{\text{out}}}{\sin \theta_f} (1 + \ln \dot{m}) \right]^{1/4} r^{-1/2}, & r_{\text{sp}} \leq r \leq r_{\text{ph}} \end{cases}$$

- $T_{\text{in}}^4 = \frac{L_{\text{Edd}} \sin(\theta_f)}{4\pi\sigma R_{\text{in}}^2}$ ,  $R_{\text{in}} = 3R_{\text{Sch}}$  and  $\sigma$  is the Stefan-Boltzmann constant.
- $\theta_f = 45^\circ$  is the funnel angle.
- $f_{\text{out}} = 0.03$  is the fraction of bolometric flux thermalised in the disc.
- $\dot{m} = \dot{M}/\dot{M}_{\text{Edd}}$ .

$$L_\nu = 2\pi R_{\text{in}}^2 \int_1^{r_{\text{ph}}} B_\nu(T(r)) r dr .$$

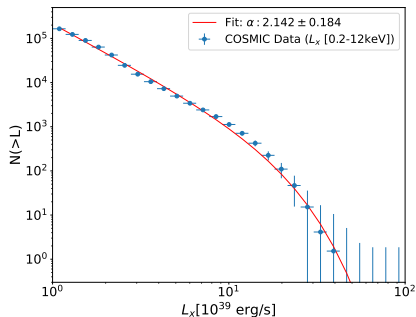
ULX population:  $L_X > 1.26 M_c \times 10^{38}$  erg/s using 0.2-12 keV band.

# Results: ULX Phase



Metal-poor massive stars are more likely to undergo direct collapse into BHs.

The power-law index  $\alpha$  varies with metallicity: Low-metallicity environments produce more high-luminosity ULXs.





## A synthetic population of ultra-luminous X-ray sources

### Optical–X-ray correlation

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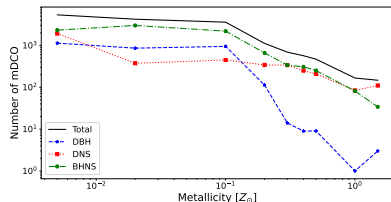
#### ABSTRACT

This paper presents an analysis of the predicted optical-to-X-ray spectral index ( $\alpha_{\text{ox}}$ ) within the context of ultra-luminous X-ray sources (ULXs) associated with stellar-mass black holes (BHs) and neutron stars (NSs). We used the population synthesis code COSMIC to simulate the evolution of binary systems and investigate the relationship between ultraviolet (UV) and X-ray emission during the ULX phase, namely the  $\alpha_{\text{ox}}$  relation. Furthermore, we investigated the impact of metallicity on  $\alpha_{\text{ox}}$  values. Notably, it predicts a significant anti-correlation between  $\alpha_{\text{ox}}$  and UV luminosity ( $L_{\text{UV}}$ ), consistent with observations. The slope of this relationship varies with metallicity for black hole ULXs (BH-ULXs). The neutron star ULX (NS-ULX) population shows a relatively consistent slope around  $-0.33$  across metallicities, with minor variations. The number of ULXs decreases with increasing metallicity, consistent with observational data. The X-ray luminosity function (XLF) shows a slight variation in its slope with metallicity, exhibiting a relative excess of high-luminosity ULXs at lower metallicities. The inclusion of the beaming effect in the analysis shows a significant impact on the XLF and  $\alpha_{\text{ox}}$ , particularly at high accretion rates, where the emission is focused into narrower cones. We found that UV emission in ULXs is predominantly disc-dominated, which is the likely origin of the  $\alpha_{\text{ox}}$  relation, with the percentage of disc-dominated ULXs increasing as metallicity rises.

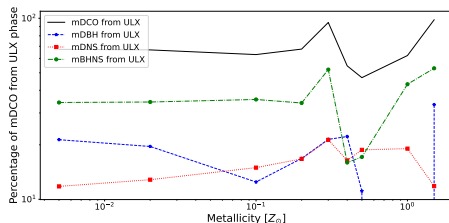
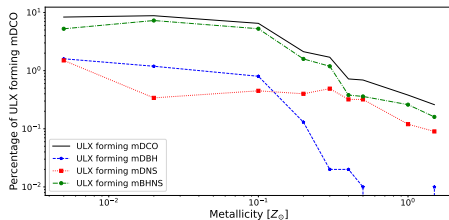
**Key words.** binaries: general

# Results: Post-ULX Compact Binary Properties

Higher metallicity increases stellar wind mass loss, reducing the number of massive stars and leading to lower DCO formation rates.



Less than 10% of ULXs evolve to mDCOs, but a striking 70 – 97% of mDCOs trace their evolutionary history through a ULX phase.



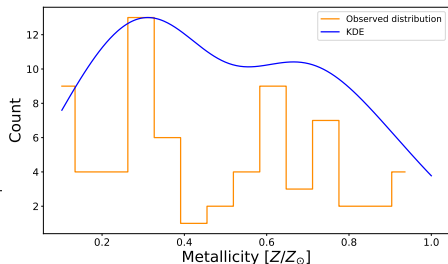
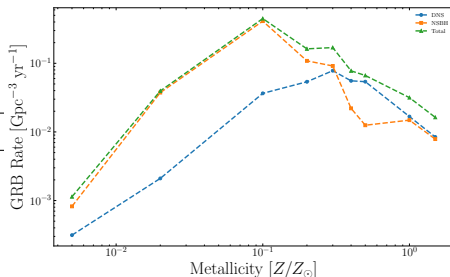
# Results: GRB rate Predictions

$$\mathcal{R}_{\text{SGRB}} = \int f_b(\theta_{\text{jet}}) P(\theta_{\text{jet}}) d\theta_{\text{jet}} \times (f_{\text{jet,DNS}} \mathcal{R}_{\text{DNS}} + f_{\text{jet,BHNS}} \mathcal{R}_{\text{BHNS}})$$

- $P(\theta_{\text{jet}})$  is the log-normal probability density function.
- $f_{\text{jet,DNS}} = \frac{N_{\text{DNS} \rightarrow \text{BH}}}{N_{\text{DNS,total}}}$
- $f_{\text{jet,BHNS}} = \frac{N_{\text{BHNS}}:(q>0.2)}{N_{\text{BHNS,total}}}$
- $f_b(\theta_{\text{jet}}) = 1 - \cos \theta_{\text{jet}}$ .
- $\mathcal{R}_{\text{BHNS}}$  and  $\mathcal{R}_{\text{DNS}}$  are BHNS and DNS merger rates taking into account the SFR and metallicity evolution with read-shift as outlined in [Bavera et al. \(2020\)](#)

# Results: GRB rate Predictions

Metallicity [ $Z/Z_{\odot}$ ]	$f_{\text{jet,BHNS}}$	$f_{\text{jet,DNS}}$
0.5%	0.726	1.000
2%	0.803	1.000
10%	0.906	0.970
20%	0.471	1.000
30%	0.449	0.972
40%	0.408	1.000
50%	0.279	1.000
100%	0.429	1.000
150%	0.391	1.000



# Torus Mass and GRB Energy from DNS Mergers

## 1. Torus Mass Estimation:

$$M_{\text{torus}} = [c_1(1 - q) + c_2] [c_3(1 + q) - M_{\text{tot}}/M_{\text{max}}]$$

- $q = M_2/M_1 \leq 1$  (mass ratio)
- $M_{\text{tot}} = M_1 + M_2$  (total mass)
- $M_{\text{max}} = \text{maximum non-rotating NS mass } (2.2 M_{\odot})$
- $c_1 = 2.974, \quad c_2 = 0.11851, \quad c_3 = 1.1193$

(Rezzolla et al. 2010) updated by Giacomazzo et al. (2013)

## 2. GRB Energy from the Torus Mass:

$$E_{\gamma, \text{iso}} = \epsilon_{\text{jet}} \cdot \epsilon_{\gamma} \cdot M_{\text{torus}} \cdot c^2$$

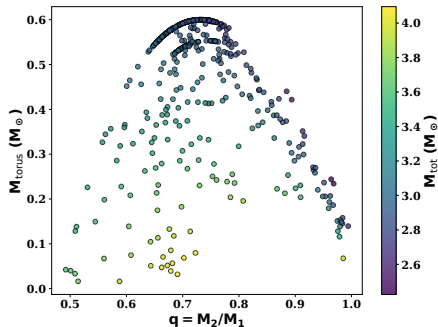
### Assumptions:

- $\epsilon_{\text{jet}}$  : Torus mass into jet energy efficiency (10%)
- $\epsilon_{\gamma}$  : jet energy into gamma rays efficiency ( 50%)

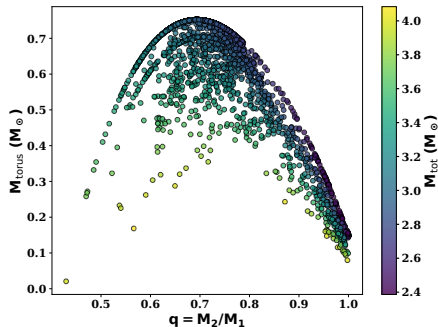
Why do this: In COSMIC  $M_f = M_1 + M_2$

# Results: Torus mass

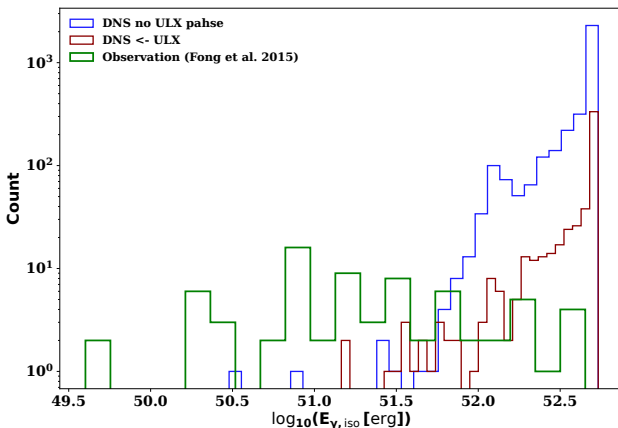
## ULX Progenitor State



## No ULX Progenitor State



# Results: Gamma-Ray Energy Distribution



- Simulated populations show a similar energy range.

# Summary and Limitations

## Summary:

- 70-97 % of mDCO comes from the ULX phase, but only 10-20 % of mDNS comes from the ULX phase.
- SGRB metallicity distribution is comparable to the observed data.
- Gamma ray energy range is slightly higher than the observed range.

## Limitations:

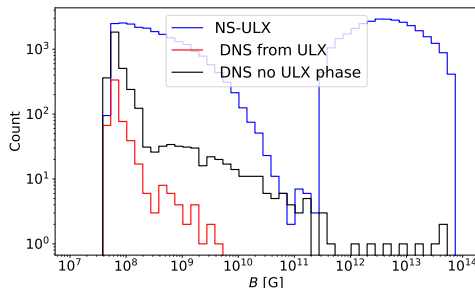
- Assumption 1:  $f_{\text{jet,DNS}} = \frac{N_{\text{DNS} \rightarrow \text{BH}}}{N_{\text{DNS,total}}}$ 
  - Jet launching is expected when the remnant mass satisfies  $M_{\text{rem}} \gtrsim 1.2 M_{\text{TOV}} \rightarrow$  hypermassive NS, collapsing to a BH after a short delay.



# Summary and Limitations

## Limitations:

- Assumption 2: No condition on magnetic field.
  - Initial magnetic field  $\sim 10^{12}$  G to produce poloidal field of  $\sim 10^{15}$  G ([Rezzolla et al. 2010](#)).
  - All our systems have  $B < 10^{12}$  G



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# Thank You!

